

A new look at the Anomalous X-ray Pulsars (AXPs)

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Basic Properties of AXPs and SGRs

- **ISOLATED** X-ray pulsars (several in **SNRs**);
- **REGULAR** spin-down at $\dot{\nu} \sim 10^{-11} - 10^{-14} \text{ Hz s}^{-1}$;
- X-ray luminosity **EXCEEDS** the spin-down power ($L_X > L_{\text{sd}} = I_{\text{ns}} \Omega \dot{\Omega}$);
- X-ray spectrum is **SOFT** ($kT \sim 0.4 - 1 \text{ keV}$);
- **CLUSTERING** of spin periods around $2 - 20 \text{ s}$;
- **Bursting activity in soft gamma-rays** ($L_\gamma \gg L_{\text{Edd}}$).

Approaches to explain AXP phenomenon

- An **unusual** Neutron Star (NS) in an **ordinary** environment
 - **Nuclear fission** (Bisnovatyi-Kogan & Chechetkin **1974**);
 - **Thermonuclear fusion** (Woosley & Wallace **1982**);
 - **Magnetic power** (Katz **1982**; Duncan & Thompson **1992**);
- An **ordinary** NS in an **unusual** environment
 - A **Low-Mass X-ray Binary powered by accretion** (Mereghetti & Stella **1995**);
 - An **isolated** NS **accreting** from a fossil disk (van Paradijs et al. **1995**);
- An **unusual** NS in an **unusual** environment
 - **Nuclear power + Magnetic power + Accretion power.**
- A **transition** of NS between **ordinary** and **unusual** states.

What powers the X-rays?

Spin-powered pulsar mechanism is excluded

$$L_X > L_{sd} = I_{ns} \Omega \dot{\Omega}$$

Residual cooling mechanism is excluded

$$L_X \sim 10^{33} - 10^{35} \text{ erg/s} \gg L_{th}$$

Additional local heating

External vs. Internal

THE VERY LOW MASS X-RAY BINARY PULSARS: A NEW CLASS OF SOURCES?

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ABSTRACT

While the distribution of spin periods of high-mass X-ray binaries spans more than four orders of magnitude (69 ms–25 minutes) the few known X-ray pulsars accreting from very low mass companions ($< 1 M_{\odot}$) have very similar periods between 5.4 and 8.7 s. These pulsars also display several other similarities, and we propose that they are members of a subclass of low-mass X-ray binaries (LMXBs) with similar magnetic field (a few times 10^{11} G), companion stars and, possibly, evolutionary histories. If they are rotating at, or close to, the equilibrium period, their properties are consistent with luminosities of the order of a few times 10^{35} ergs s^{-1} . These pulsars might represent the closest members of a subclass of LMXBs characterized by lower luminosities, higher magnetic fields, and smaller ages than nonpulsating LMXBs.

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On the nature of the ‘anomalous’ 6-s X-ray pulsars

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Abstract. Recently it has become clear that there is a group of X-ray pulsars, with pulse periods close to 6 seconds which are distinct from accreting strongly magnetized neutron stars in X-ray binary systems. Here we argue that these objects are the recent products of the evolution of massive stars. They are unlikely to be neutron stars that formed through accretion induced collapse of a white dwarf. We propose that they are single neutron stars accreting from a disk, the recent remnants of the common-envelope evolution of a high-mass X-ray binary.

State by 1995

- AXPs in X-rays resemble Accretion-powered Pulsars
 - a Neutron Star (NS) in a LMXB, or
 - a NS accreting from a fossil Keplerian disk
- AXPs rotate at the equilibrium period
(for a Keplerian disk $P_{\text{eq}}^{(\text{Kd})} \sim \text{a few seconds}$)

1998: Magnetars, magnetars and only magnetars!..

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An X-ray pulsar with a superstrong magnetic field in the soft γ -ray repeater SGR1806 – 20

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Here we report the discovery of pulsations in the persistent X-ray flux of SGR1806 – 20, with a period of 7.47 s and a spindown rate of $2.6 \times 10^{-3} \text{ s yr}^{-1}$. We argue that the spindown is due to magnetic dipole emission and find that the pulsar age and (dipolar) magnetic field strength are $\sim 1,500$ years and 8×10^{14} gauss, respectively. Our observations demonstrate the existence of ‘magnetars’, neutron stars with magnetic fields about 100 times stronger than those of radio pulsars,

Spin-down to the period of 7.47 s on a timescale of only 1,500 years!

Regularly spinning-down like a radiopulsar

$$(P_{\text{eq}} \gg P_{\text{s}})$$

The spin-down torque inferred from observations is larger than the maximum possible torque expected in accretion scenarios

But why do AXPs resemble Accretion-powered Pulsars?

Can SGR 1806-20 be an accretion-powered pulsar?

Accretion process: $r_m < r_{\text{cor}} = \left(\frac{GM_{\text{ns}}}{\omega_s^2} \right)^{1/3} \simeq 6.6 \times 10^8 \text{ cm}$

Regular spin-down: $|K_{\text{sd}}| \geq 2\pi I |\dot{\nu}_{\text{sd}}^{\text{obs}}|$

- Quasi-spherical accretion

$$|K_{\text{sd}}^{(\text{sp})}| \leq k_t \dot{m} \omega_s r_A^2$$

$$\implies \underline{r_A \geq 12 r_{\text{cor}}} \times k_t^{-1/2}$$

- Keplerian disk accretion

$$|K_{\text{sd}}^{(\text{Kd})}| \leq k_t \mu^2 / r_{\text{cor}}^3$$

$$\implies \underline{r_A \geq 4 r_{\text{cor}}} \times k_t^{-2/7}$$

By **1998** the answer was **negative**...

This indicates that SGR 1806-20 is

- either a non-accreting NS,

- or the accretion onto this NS differs from both
 - Quasi-spherical, and
 - Keplerian disk

Magnetic-Levitation Accretion onto a Neutron Star

1. **Accretion** from a **magnetized** wind ($\beta_0 \sim 1$)
2. **Deceleration** of the free-falling material at the **Shvartsman radius** R_{sh}
3. **Formation** of the **non-Keplerian Magnetically-Levitating Disk (MAGLEV Disk)**
4. **Diffusion** of **accreting material** into the **stellar MF** at the magnetospheric boundary

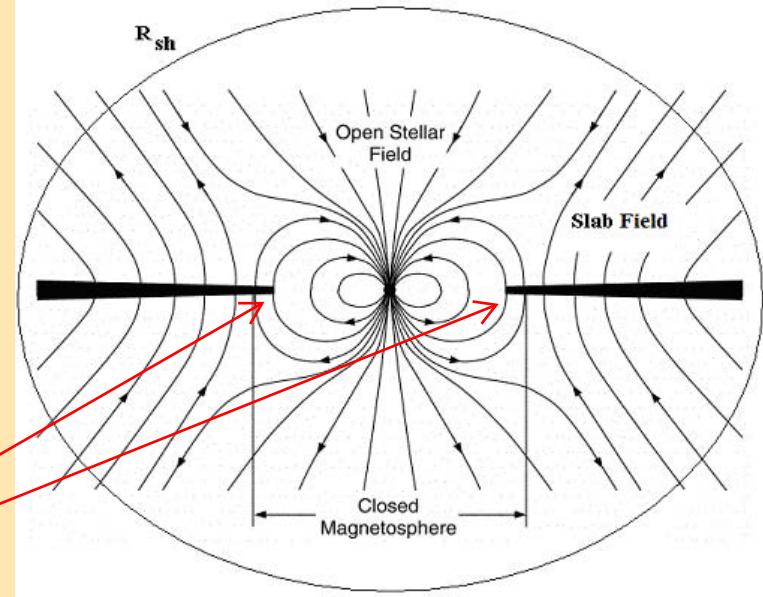
Shvartsman radius	$R_{\text{sh}} = \beta_0^{-2/3} \left(\frac{c_s(R_G)}{v_{\text{rel}}} \right)^{4/3} R_G$
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New parameters:	Magnetospheric radius	$r_{\text{ma}} = \left(\frac{c m_p^2}{16 \sqrt{2} e k_B} \right)^{2/13} \frac{\alpha_B^{2/13} \mu^{6/13} (GM_{\text{ns}})^{5/13}}{T_0^{2/13} L_x^{4/13} R_{\text{ns}}^{4/13}}$
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ML-torque	$K_{\text{ml}} = k_m \frac{\mu^2}{(r_{\text{ma}} r_{\text{cor}})^{3/2}} \left(1 - \frac{\Omega_{\text{ml}}(r_{\text{ma}})}{\Omega_s} \right)$
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Magnetically-Levitating (MAGLEV) Disk

$$\left\{ \begin{array}{l} \frac{\mu^2}{2\pi r_{\text{ma}}^6} = \rho(r_{\text{ma}}) c_s^2(r_{\text{ma}}) \\ \dot{M}_{\text{in}}(r_{\text{ma}}) = \frac{L_X R_{\text{ns}}}{GM_{\text{ns}}} \\ \dot{M}_{\text{in}}(r_{\text{ma}}) = 4\pi r_{\text{ma}} \delta_m \rho(r_{\text{ma}}) v_{\text{ff}}(r_{\text{ma}}) \\ \delta_m(r_{\text{ma}}) = \left[t_{\text{ff}}(r_{\text{ma}}) D_{\text{eff}}(r_{\text{ma}}) \right]^{1/2} \end{array} \right.$$



$$D_{\text{eff}}(r_{\text{ma}}) = \alpha D_B(r_{\text{ma}}) = \alpha \frac{ck_B T_i(r_{\text{ma}})}{16eB(r_{\text{ma}})}$$

$$r_{\text{ma}} = \left(\frac{cm_p^2}{16\sqrt{2}ek_B} \right)^{2/13} \frac{\alpha^{2/13} \mu^{6/13} (GM_{\text{ns}})^{5/13}}{T_0^{2/13} L_X^{4/13} R_{\text{ns}}^{4/13}}$$

Magnetic field of **AXPs** and **SGRs** in **MLA** scenario

- NS in the accretor state $r_{\text{ma}} < r_{\text{cor}}$ (An upper limit)

$$\mu < 10^{31} \text{ G cm}^3 \times \alpha_{0.1}^{-1} T_6^{1/3} m^{-1/9} L_{34}^{2/3} R_6^{2/3} \left(\frac{P_s}{5 \text{ s}} \right)^{13/9}$$

- Spin-down rate $|K_{\text{ml}}| \geq 2\pi I |\dot{\nu}_{\text{sd}}|$ (A lower limit)

$$\mu \geq \left[\frac{2\pi I |\dot{\nu}_{\text{sd}}|}{k_t} \right]^{1/2} \left(r_{\text{ma}} r_{\text{cor}} \right)^{3/4}$$

Parameters of SGRs and AXPs in the MLA scenario

Name	$P_s,$ c	$ \dot{\nu}_{sd} ,$ $10^{-12} \text{ Hz s}^{-1}$	$L_X,$ $10^{34} \text{ erg s}^{-1}$	$B_*,$ 10^{12} G	$r_{\text{cor}},$ 10^8 cm	$r_{\text{ma}},$ 10^8 cm	$a_p,$ 10^5 cm	$T_{\text{bb}},$ keV
SGR 1627-41	2.6	2.8	0.25	3.9	3.2	1.2	0.9	0.5
SGR 1900+14	5.2	3.4	9.0	2.7	5.2	1.0	1.0	1.1
SGR 1806-20	7.6	13	16	12	6.6	1.7	0.8	1.4
SGR 0526-66	8.05	0.59	14	1.2	6.9	0.6	1.3	1.1
1E 1547.0-5408	2.07	3.2	14	1.5	2.8	1.3	0.9	1.3
CXOU J174505.7-381031	3.83	4.4	6.0	2.3	4.2	1.0	1.0	0.98
PSR J1622-4950	4.33	0.9	0.063	1.5	4.6	1.2	0.9	0.33
XTE J1810-197	5.54	0.26	3.9	0.3	5.4	0.46	1.5	0.73
1E 1048.1-5937	6.45	0.55	0.6	0.31	5.95	0.83	1.1	0.53
1E 2259+586	6.98	0.01	2.2	0.024	6.3	0.17	2.4	0.49
CXOU J010043.1-721134	8.02	0.29	6.1	0.51	6.9	0.51	1.4	0.83
4U 0142+61	8.69	0.03	11	0.11	7.26	0.21	2.2	0.77
CXO J164710.2-455216	10.61	0.006	0.3	0.011	8.3	0.22	2.1	0.32
1RXS J170849.0-400910	11.0	0.16	5.9	0.4	8.5	0.46	1.5	0.81
1E J1841-045	11.8	0.28	19	0.99	8.9	0.49	1.4	1.1

- **MAGNETIC FIELD:** $B_* \sim (0.01 - 10) \times 10^{12} \text{ G}$

- **SOFT X-ray SPECTRUM** $kT \sim 0.3 - 1.4 \text{ keV}$

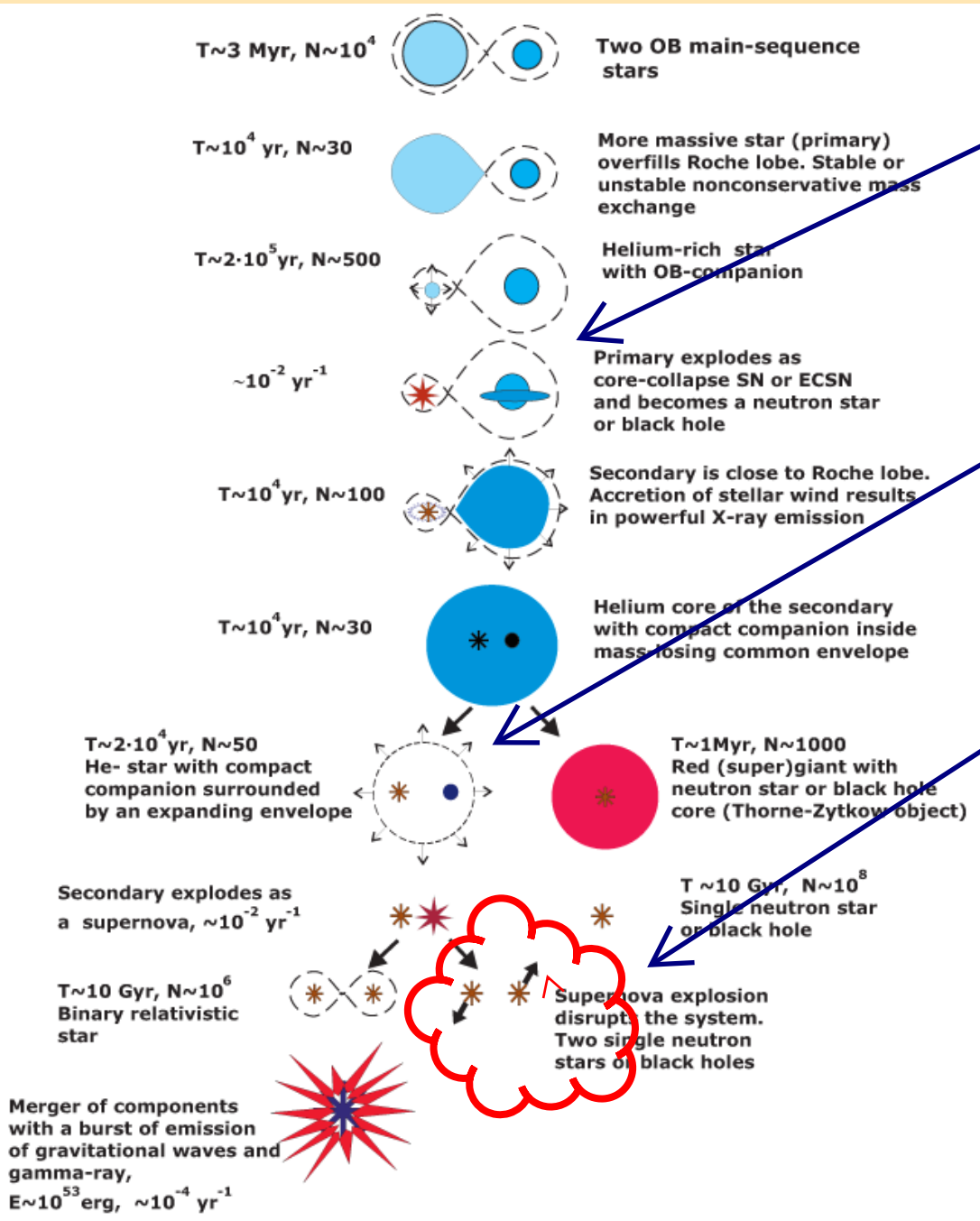
How could a moderately magnetized NS spin-down to a period of a few seconds on a timescale of only 1000 yr?

Ejector spin-down time $\tau_{\text{ej}} \sim 10^6 \text{ yr} \times \mu_{30}^{-1} I_{45} \mathfrak{M}_{15}^{-1/2} v_8^{-1/2}$

Can AXP's be older than $\sim 10^6$ years?

- The spin-down timescale is $\tau_{\text{sd}} \sim P_{\text{s}}/2\dot{P} \sim 200 - 10^5 \text{ yr}$
 - Was the NS born $\sim \tau_{\text{sd}}$ years ago, or
 - it had switched into a new (current) state $\sim \tau_{\text{sd}}$ years ago?
- The Age of SNR is $\tau_{\text{SNR}} \sim 10^3 - 10^4 \text{ yr}$
 - Was the NS born about τ_{SNR} years ago, or
 - it has nothing to do with this SNR?

AXPs as DESCENDANTS of High Mass X-ray Binaries



1. A NS forms in a Massive Binary System during the **1st SN explosion**

2. Its spin period evolves on $\sim 10^6 - 10^7 \text{ yr}$ to

$$P_{\text{eq}}^{(\text{Kd})} \simeq 3 \text{ s} \times \mu_{30}^{6/7} m^{5/7} \dot{M}_{17}^{-3/7}$$

3. After **2nd SN explosion** the system is **disrupted** the old ($\sim 10^6 - 10^7 \text{ yr}$) NS becomes **isolated**

– it is **embedded into a SNR**

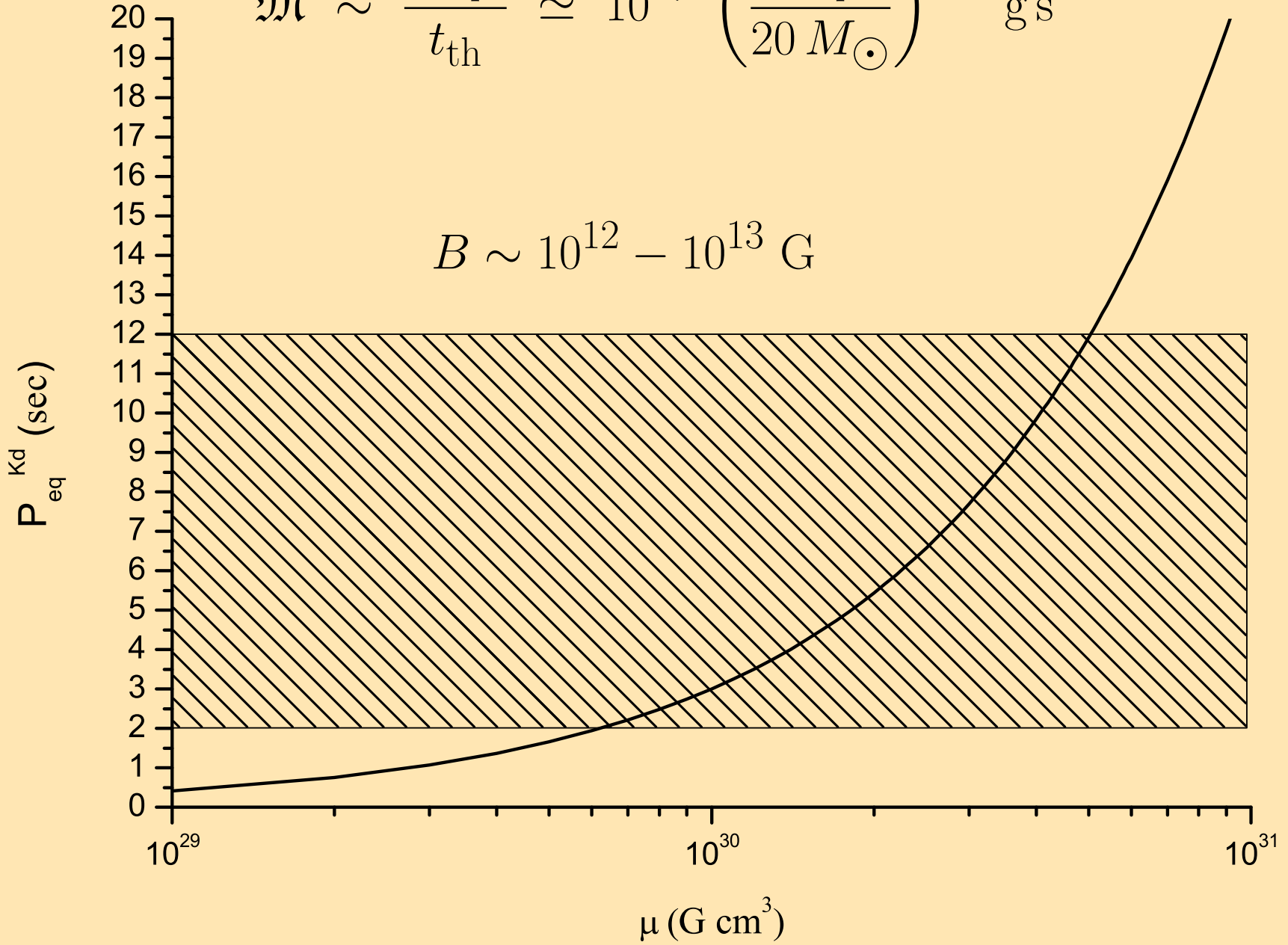
– it **captures material** from SN ejecta

– it **rotates** with the **period** $P_{\text{eq}}^{(\text{Kd})}$

– it is **spinning-down** on $200 - 10^5 \text{ yr}$

$$\dot{m} \sim \frac{M_{\text{opt}}}{t_{\text{th}}} \simeq 10^{17} \left(\frac{M_{\text{opt}}}{20 M_{\odot}} \right)^{-1} \text{ g s}^{-1}$$

$$B \sim 10^{12} - 10^{13} \text{ G}$$



$$t_{\text{th}} = GM_{\text{opt}}^2 / (R_{\text{opt}} L_{\text{opt}}) \sim 3 \times 10^7 \text{ yr}$$

Mass of the fossil ML-disk is

$$M_d = 4\pi \int_{r_{\text{ma}}}^{R_{\text{sh}}} \rho(r) h_z(r) r dr$$

Disk parameters

- Temperature $T(r) = \left(\frac{\dot{m} G M_{\text{ns}}}{4\pi r^3 \sigma_{\text{SB}}} \right)^{1/4}$; $T(r) \propto r^{-3/4}$
- Halfthickness $h_z(r) = \left(\frac{k_B T(r) r^3}{m_p G M_{\text{ns}}} \right)^{1/2}$; $h_z(r) \propto r^{-9/8}$
- Density $\rho(r) c_s^2(r) \propto r^{-5/2}$; $\rho(r) \propto r^{-7/4}$

$$M_d \simeq 10^{-7} M_{\odot} \times \mu_{30}^{5/13} \dot{m}_{20}^{99/104} c_7^{11/6} v_8^{-55/12} \times \left[\alpha_{0.1}^{-7/3} \beta_0^{-11/12} m^{25/52} \right]$$

The ML-disk can be formed from the innermost slowly moving part of the SN ejecta

CONCLUSIONS

AXPs are descendants of HMXBs accreting from a residual MAGLEV-Disk

- Neutron Stars rotating with the period $P_{\text{eq}}^{\text{Kd}} \sim$ a few seconds;
- Isolated and rather old ($\sim 10^6 - 10^7$ yr)
- Regular magnetic field ($B_* \sim (1 - 10) \times 10^{12}$ G)

What powers activity of SGRs and AXPs in gamma-rays?

- *The energy source is located inside the Neutron Star*
 - Magnetic energy in the crust is $E_m \leq 5 \times 10^{46} B_{15}^2$ erg;
 - Binding energy in the crust is $E_{\text{nuc}} \sim 10^{47}$ erg (for $\rho \sim (0.4 - 10) \times 10^{11}$ g cm⁻³)
- *Questions about the trigger and transmission are still open*